

# Series lectures of phase-field model

## 03. Cahn-Hilliard Equation I

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- 1 Motivation of phase-field model
- 2 Cahn-Hilliard Equation
  - Derivation of Cahn-Hilliard Equation



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# Growth of a precipitate

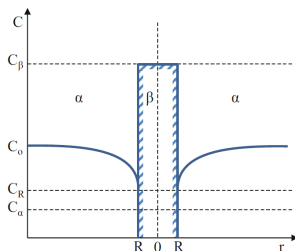
- Diffusion equation (Fick's 2nd law)

$$\frac{\partial c(\mathbf{r}, t)}{\partial t} = D \nabla^2 c(\mathbf{r}, t)$$

- Mass balance at the interface

$$\frac{dR}{dt} (c^\beta - c^\alpha) = D \left. \frac{\partial c(\mathbf{r}, t)}{\partial r} \right|_{r=R(t)}$$

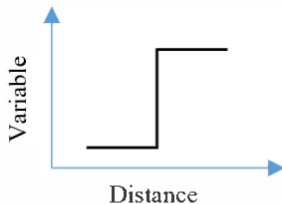
- Free boundary problems



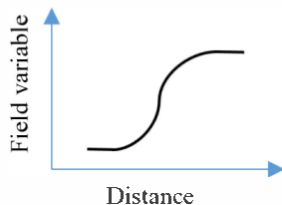
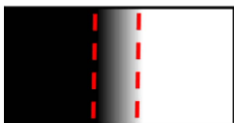
# Sharp interface approach

Write one (or many) PDE that holds everywhere, in both phases and at the interface. **Do not need to track the location of the boundary.**

Sharp interface



Diffuse interface



- Conserved order parameter: Cahn-Hilliard equation
- Non-conserved order parameter: Allen-Cahn equation

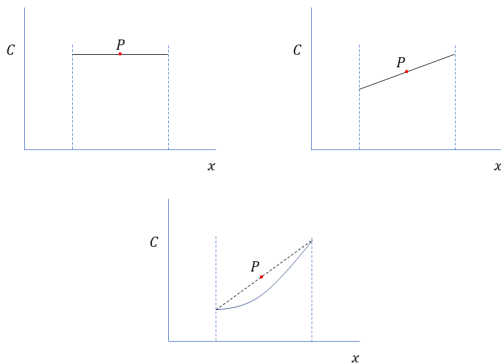
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# Derivation of Cahn-Hilliard equation

- Cahn-Hilliard (JCP, 1958): Consider the energy associated with gradients in composition.





# Derivation of Cahn-Hilliard equation

- The chemical potential is function of  $\frac{d^2c}{dx^2}$

$$f = f\left(c, \frac{dc}{dx}, \frac{d^2c}{dx^2}, \dots\right)$$

- By Taylor's expansion with omitting higher-order terms,

$$f(c, dc/dx) = f(c, 0) + L \frac{dc}{dx} + \kappa \left(\frac{dc}{dx}\right)^2$$

where

$$L = \frac{\partial f}{\partial (dc/dx)} \quad \kappa = \frac{1}{2} \frac{\partial f}{\partial ((dc/dx)^2)}$$

- $L = 0$  for centrosymmetric crystal,  $f(c)$  is free energy per volume.

$$F = A \int_V \left[ f(c) + \kappa \left(\frac{dc}{dx}\right)^2 \right] dV$$



# Derivation of Cahn-Hilliard equation

- For 1-D case

$$F = A \int_{x_1}^{x_2} \left[ f(c) + \kappa \left( \frac{dc}{dx} \right)^2 \right] dx$$
$$\delta F = A \int_{x_1}^{x_2} \left[ \frac{\partial f(c)}{\partial c} \delta c + \kappa \delta \left( \frac{dc}{dx} \right)^2 \right] dx$$
$$= A \int_{x_1}^{x_2} \left[ \frac{\partial f(c)}{\partial c} \delta c + 2\kappa \left( \frac{dc}{dx} \right) \delta \left( \frac{dc}{dx} \right) \right] dx$$

- By integrating by parts

$$\int_{x_1}^{x_2} \left[ \left( \frac{dc}{dx} \right) \delta \left( \frac{dc}{dx} \right) \right] dx = \frac{dc}{dx} \delta c \Big|_{x_1}^{x_2} - \int_{x_1}^{x_2} \frac{d^2c}{dx^2} \delta c dx$$

when  $\frac{dc}{dx}$  is 0 when  $x = x_1, x_2$  then

$$\int_{x_1}^{x_2} 2\kappa \left( \frac{dc}{dx} \right) \delta \left( \frac{dc}{dx} \right) dx = -2\kappa \int_{x_1}^{x_2} \frac{d^2c}{dx^2} \delta c dx$$



# Derivation of Cahn-Hilliard equation

- At equilibrium,

$$\frac{\delta F}{\delta c} = A \int_{x_1}^{x_2} \left[ \frac{\partial f(c)}{\partial c} - 2\kappa \frac{d^2 c}{dx^2} \right] dx = 0$$

therefore,

$$\frac{\partial f(c)}{\partial c} - 2\kappa \frac{d^2 c}{dx^2} = 0$$

- For the case of non-conserved order parameter( $\eta$ ), means no constrain on the average value of  $\eta$ , we can state

$$\frac{\partial f(\eta)}{\partial \eta} - 2\kappa \frac{d^2 \eta}{dx^2} = 0$$

- For conserved order parameter, such as composition( $c$ ), it have to be conserved

$$\int_{x_1}^{x_2} [c(x) - c_0] dx = 0$$

where  $c_0$  is the nominal alloy composition.



# Derivation of Cahn-Hilliard equation

- Then we have

$$F = A \int_{x_1}^{x_2} \left[ f(c) + \kappa \left( \frac{dc}{dx} \right)^2 - \lambda (c(x) - c_0) \right] dx$$

- Take variation

$$\begin{aligned} \delta F &= A \int_{x_1}^{x_2} \left[ \delta \left[ f(c) + \kappa \left( \frac{dc}{dx} \right)^2 \right] - \lambda \delta (c(x) - c_0) \right] dx \\ &= A \int_{x_1}^{x_2} \left[ \frac{\partial f}{\partial c} - 2\kappa \frac{d^2 c}{dx^2} - \lambda \right] \delta c dx \end{aligned}$$

therefore, we reach the Euler-Lagrange equation

$$\lambda = \frac{\partial f}{\partial c} - 2\kappa \frac{d^2 c}{dx^2}$$

- $\lambda$  have to be uniform at equilibrium.



# Derivation of Cahn-Hilliard equation

- If  $\lambda$  is not uniform, flux have to be present. Therefore, we can write

$$J_A = -M\nabla\lambda = -M\nabla\left(\frac{\partial f}{\partial c_A} - 2\kappa\frac{\partial^2 c_A}{\partial x^2}\right)$$

where

$$f = c_0[\mu_A c_A + \mu_B c_B]$$



- Applying mass conservation equation

$$\begin{aligned}\frac{\partial c}{\partial t} &= -\frac{\partial J}{\partial x} \\ &= \frac{d}{dx} \cdot \left[ M \frac{\partial}{\partial x} \left( \frac{\partial f}{\partial c} - 2\kappa \frac{\partial^2 c}{\partial x^2} \right) \right]\end{aligned}$$

we reach the Cahn-Hilliard equation in one-dimensional system.

# Derivation of Cahn-Hilliard equation

- Assume constant mobility,

$$\begin{aligned}\frac{\partial c}{\partial t} &= M \frac{\partial}{\partial x} \left[ \frac{\partial}{\partial x} \left( \frac{\partial f}{\partial c} - 2\kappa \frac{\partial^2 c}{\partial x^2} \right) \right] \\ &= M \frac{\partial}{\partial x} \left[ \frac{\partial^2 f}{\partial c^2} \frac{\partial c}{\partial x} - 2\kappa \frac{\partial^3 c}{\partial x^3} \right] \\ &= M \left[ \frac{\partial}{\partial x} \left( \frac{\partial^2 f}{\partial c^2} \frac{\partial c}{\partial x} \right) - 2\kappa \frac{\partial^4 c}{\partial x^4} \right]\end{aligned}$$

- Unfortunately, closed form of the solution does not exist generally.

