

Series lectures of phase-field model

01. Quick Review of Thermodynamics

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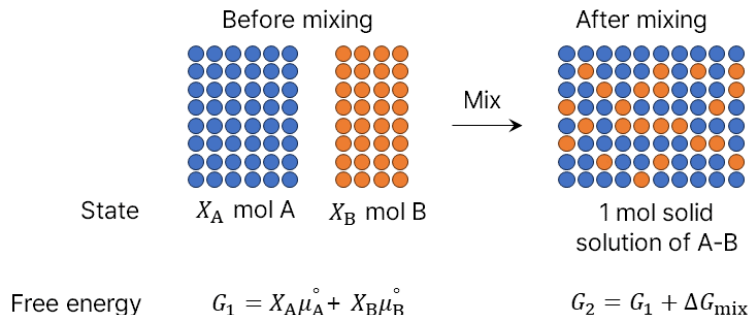


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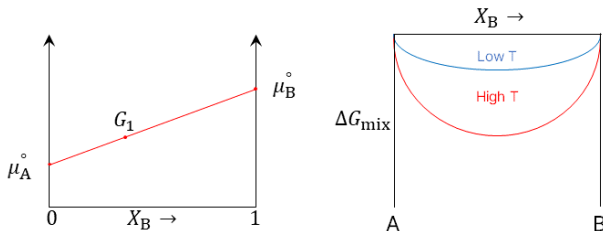
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 - Thermodynamic equilibrium

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Quick review for solution thermodynamics



Ideal solution



The integral molar Gibbs free energy of ideal solution is

$$G_m^{\text{id}} = X_A \mu_A^\circ + X_B \mu_B^\circ + RT [X_A \ln X_A + X_B \ln X_B]$$

Extend to Regular solution

	Ideal	Regular
ΔH_{mix}	0	$\Omega X_A X_B$
ΔS_{mix}	$-R [X_A \ln X_A + X_B \ln X_B]$	$-R [X_A \ln X_A + X_B \ln X_B]$
ΔG_{mix}	$RT [X_A \ln X_A + X_B \ln X_B]$	$\Omega X_A X_B + RT [X_A \ln X_A + X_B \ln X_B]$
ΔV_{mix}	0	0

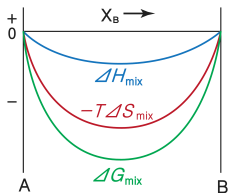
Table: Comparison between Ideal and Regular solution

The integral molar Gibbs free energy of regular solution is

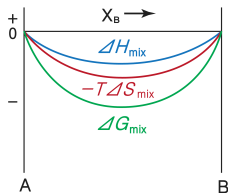
$$G_m^{\text{reg}} = X_A \mu_A^\circ + X_B \mu_B^\circ + RT [X_A \ln X_A + X_B \ln X_B] + \Omega X_A X_B$$



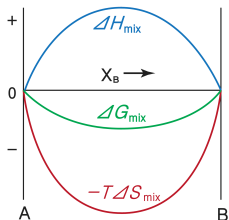
Behavior of regular solution



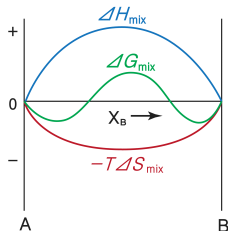
(a) $\Omega < 0$, high T



(b) $\Omega < 0$, low T



(c) $\Omega > 0$, high T



(d) $\Omega > 0$, low T

2nd law of thermodynamics in the Materials science

- The system evolves to minimize the Gibbs energy of the system.
- Equilibrium means that chemical potential of every species in every phase is same under constant T and p .

Integral molar quantity and Partial molar quantity

For binary solution, the chemical potential defined by the J/atom is

$$\mu_A = \left(\frac{\partial G}{\partial n_A} \right)_{n_B} \quad \mu_B = \left(\frac{\partial G}{\partial n_B} \right)_{n_A}$$

where n_A and n_B are numbers of A and B atoms.

$$\begin{aligned} \mu_A &= \left(\frac{\partial G}{\partial n_A} \right)_{n_B} = \left(\frac{\partial (nG_m)}{\partial n_A} \right)_{n_B} \\ &= G_m + n \cdot \left(\frac{\partial G_m}{\partial X_A} \right)_{X_B} \left(\frac{\partial X_A}{\partial n_A} \right)_{n_B} + n \cdot \left(\frac{\partial G_m}{\partial X_B} \right)_{X_A} \left(\frac{\partial X_B}{\partial n_A} \right)_{n_B} \end{aligned} \quad (1)$$

where

$$\begin{aligned} n &= n_A + n_B \\ X_A &= \frac{n_A}{n} \quad X_B = \frac{n_B}{n} \end{aligned}$$



Since

$$\left(\frac{\partial X_A}{\partial n_A}\right)_{n_B} = \frac{1}{n} - \frac{n_A}{n^2} = \frac{n - n_A}{n^2} = \frac{1 - X_A}{n} = \frac{X_B}{n}$$
$$\left(\frac{\partial X_B}{\partial n_A}\right)_{n_B} = -\frac{n_B}{n^2} = -\frac{X_B}{n}$$

Put it back to Eq. 1,

$$\mu_A = G_m + X_B \left[\left(\frac{\partial G_m}{\partial X_A}\right)_{X_B} - \left(\frac{\partial G_m}{\partial X_B}\right)_{X_A} \right]$$

Consistently,

$$\mu_B = G_m + X_A \left[\left(\frac{\partial G_m}{\partial X_B}\right)_{X_A} - \left(\frac{\partial G_m}{\partial X_A}\right)_{X_B} \right]$$

Then we can introduce the diffusion potential $\mu_A - \mu_B$ by

$$\mu_A - \mu_B = \left(\frac{\partial G_m}{\partial X_A}\right)_{X_B} - \left(\frac{\partial G_m}{\partial X_B}\right)_{X_A}$$



The total derivative is

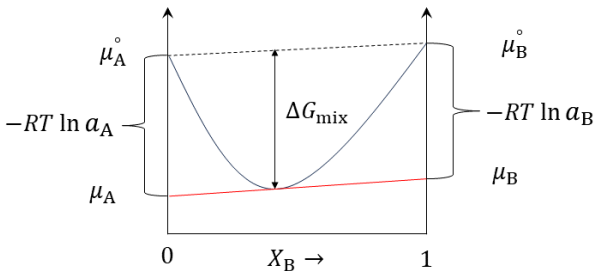
$$\frac{dG_m}{dX_A} = \left(\frac{\partial G_m}{\partial X_A} \right)_{X_B} \frac{\partial X_A}{\partial X_A} + \left(\frac{\partial G_m}{\partial X_B} \right)_{X_A} \frac{\partial X_B}{\partial X_A} = \left(\frac{\partial G_m}{\partial X_A} \right)_{X_B} - \left(\frac{\partial G_m}{\partial X_B} \right)_{X_A}$$

Consistently,

$$\frac{dG_m}{dX_B} = \left(\frac{\partial G_m}{\partial X_B} \right)_{X_A} - \left(\frac{\partial G_m}{\partial X_A} \right)_{X_B} \quad (2)$$

therefore,

$$\mu_B = G_m - X_A \frac{dG_m}{dX_A} \quad \mu_A = G_m - X_B \frac{dG_m}{dX_B} \quad (3)$$



From Eq.2, therefore, we can verify that the y intercept of tangential line is the chemical potential of two elements.

$$G_m(X_B = 0) = G_m - X_B \times \left(\frac{dG_m}{dX_B} \right) = \mu_A$$

Chemical potential of regular solution

For regular solution,

$$G_m = X_A \mu_A^\circ + X_B \mu_B^\circ + RT [X_A \ln X_A + X_B \ln X_B] + \Omega X_A X_B$$

From Eq. 3,

$$\begin{aligned}\mu_B &= X_A \mu_A^\circ + X_B \mu_B^\circ + RT [X_A \ln X_A + X_B \ln X_B] + \Omega X_A X_B \\ &\quad - X_A \left[\mu_A^\circ - \mu_B^\circ + RT [\ln X_A - \ln X_B] + \Omega X_B - \Omega X_A \right] \\ &= \mu_B^\circ + RT \ln X_B + \Omega X_A^2\end{aligned}$$

Consistently,

$$\mu_A = \mu_A^\circ + \Omega(1 - X_A)^2 + RT \ln X_A$$

$$\mu_B = \mu_B^\circ + \Omega(1 - X_B)^2 + RT \ln X_B$$



Chemical equilibrium

For binary A-B system, when α - β phases are under equilibrium,

$$\mu_A^\alpha(X_A^{\alpha,e}) = \mu_A^\beta(X_A^{\beta,e}) \quad \mu_B^\alpha(X_B^{\alpha,e}) = \mu_B^\beta(X_B^{\beta,e})$$

where e superscript denotes the equilibrium state. Using Eq. 3, the first equation can be re-written,

$$G_m^\alpha(X_B^{\alpha,e}) - X_B^{\alpha,e} \left. \frac{dG_m^\alpha}{dX_B} \right|_{X_B^{\alpha,e}} = G_m^\beta(X_B^{\beta,e}) - X_B^{\beta,e} \left. \frac{dG_m^\beta}{dX_B} \right|_{X_B^{\beta,e}}$$

When

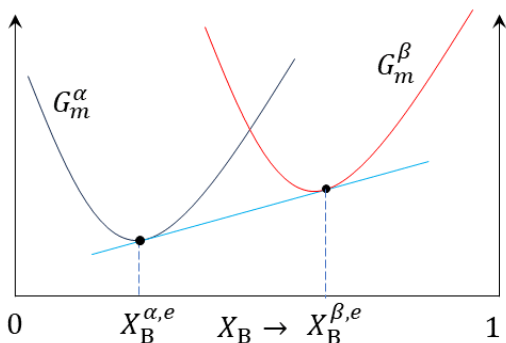
$$\left. \frac{dG_m^\alpha}{dX_B} \right|_{X_B^{\alpha,e}} = \left. \frac{dG_m^\beta}{dX_B} \right|_{X_B^{\beta,e}} = \frac{dG_m}{dX_B}$$

Proceed to

$$\frac{dG_m}{dX_B} = \frac{G_m^\alpha(X_B^{\alpha,e}) - G_m^\beta(X_B^{\beta,e})}{X_B^{\alpha,e} - X_B^{\beta,e}}$$



The condition is named as the *common tangent condition*, which means the system is under the equilibrium, which can be visualized as below:



On the common tangent line, the chemical potentials of the A, B solutes are constant, therefore, the activities are constant within two-phase region.