

Thermodynamics of materials

12. Construction of Fundamental Equation of Thermodynamics

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Table of Contents

- 1 Choice of Fundamental Equation of Thermodynamics
- 2 Molar Helmholtz Free energy change
- 3 Chemical Potential change
- 4 Chemical Potential change with phase change
- 5 Reference States
 - Reference State for Helmholtz free energy
 - Reference State for Chemical Potential
- 6 Fundamental Equations with Multiphysics Effects
 - Electrochemical Potential
 - Chemical Potential Including Surface



Table of Contents

- 1 Choice of Fundamental Equation of Thermodynamics
- 2 Molar Helmholtz Free energy change
- 3 Chemical Potential change
- 4 Chemical Potential change with phase change
- 5 Reference States
 - Reference State for Helmholtz free energy
 - Reference State for Chemical Potential
- 6 Fundamental Equations with Multiphysics Effects
 - Electrochemical Potential
 - Chemical Potential Including Surface



Choice of Fundamental Equation of Thermodynamics

- In the materials science, the commonly chosen natural variables are (T, s) and (p, v) , therefore, the most commonly employed fundamental equations in molar quantities are

$$f(T, v) = u - Ts = -pv + \mu \rightarrow df = -sdT - pdv$$

$$g(T, p) = u - Ts + pv = \mu \rightarrow dg = d\mu = -sdT + vdp$$



Table of Contents

- 1 Choice of Fundamental Equation of Thermodynamics
- 2 Molar Helmholtz Free energy change
- 3 Chemical Potential change
- 4 Chemical Potential change with phase change
- 5 Reference States
 - Reference State for Helmholtz free energy
 - Reference State for Chemical Potential
- 6 Fundamental Equations with Multiphysics Effects
 - Electrochemical Potential
 - Chemical Potential Including Surface

Molar Helmholtz Free energy change

- We have

$$f(T, v) = f^\circ(T_0, v_0) - \int_{T_0}^T s dT - \int_{v_0}^v p dv$$

with

$$f^\circ(T_0, v_0) = u^\circ(T_0, v_0) - T_0 s^\circ(T_0, v_0)$$

$$f(T, v) = u(T, v) - Ts(T, v)$$

- We have

$$u(T, v) - u^\circ(T_0, v_0) = \int_{T_0}^T c_v dT + T \int_{v_0}^v \alpha B_T dv - \int_{v_0}^v p dv$$

$$s(T, v) - s^\circ(T_0, v_0) = \int_{T_0}^T c_v d(\ln T) + \int_{v_0}^v \alpha B_T dv$$



Molar Helmholtz Free energy change

- Proceed to

$$f(T, v) - f^\circ(T_0, v_0) = u(T, v) - u^\circ(T_0, v_0) - Ts(T, v) + T_0s^\circ(T_0, v_0)$$

- RHS is

$$\int_{T_0}^T c_v dT + T \int_{v_0}^v \alpha B_T dv - \int_{v_0}^v p dv + T_0 s^\circ(T_0, v_0) - T \left[s^\circ(T_0, v_0) + \int_{T_0}^T c_v d(\ln T) + \int_{v_0}^v \alpha B_T dv \right]$$

- Therefore,

$$- \int_{T_0}^T s dT = \int_{T_0}^T c_v dT - s^\circ(T_0, v_0)(T - T_0) - T \int_{T_0}^T c_v d(\ln T)$$



Table of Contents

- 1 Choice of Fundamental Equation of Thermodynamics
- 2 Molar Helmholtz Free energy change
- 3 Chemical Potential change**
- 4 Chemical Potential change with phase change
- 5 Reference States
 - Reference State for Helmholtz free energy
 - Reference State for Chemical Potential
- 6 Fundamental Equations with Multiphysics Effects
 - Electrochemical Potential
 - Chemical Potential Including Surface



Chemical Potential change

- Since we have at reference state,

$$\mu^\circ(T_0, p_0) = h^\circ(T_0, p_0) - T_0 s^\circ(T_0, p_0)$$

and at state (T, p)

$$\mu(T, p) = h(T, p) - T s(T, p)$$

- The enthalpy and entropy are given by

$$h(T, p) = h^\circ(T_0, p_0) + \int_{T_0}^T c_p dT + \int_{p_0}^p v(1 - T\alpha) dp$$

$$s(T, p) = s^\circ(T_0, p_0) + \int_{T_0}^T \frac{c_p}{T} dT - \int_{p_0}^p v\alpha dp$$



- Then

$$\begin{aligned}\mu(T, p) = & \mu^\circ(T_0, p_0) + \int_{T_0}^T [c_p - s^\circ(T_0, p_0)] dT + \int_{p_0}^p v(1 - T\alpha) dp \\ & - T \left[\int_{T_0}^T \frac{c_p}{T} dT - \int_{p_0}^p v\alpha dp \right]\end{aligned}$$

therefore,

$$\mu(T, p) = \mu^\circ(T_0, p_0) + \int_{T_0}^T [c_p - s^\circ(T_0, p_0)] dT - T \int_{T_0}^T \frac{c_p}{T} dT + \int_{p_0}^p v dp$$

- The volume is given by

$$v(T, p) = v_0(T_0, p_0) \exp [\alpha(T - T_0) - \beta_T(p - p_0)]$$



Table of Contents

- 1 Choice of Fundamental Equation of Thermodynamics
- 2 Molar Helmholtz Free energy change
- 3 Chemical Potential change
- 4 Chemical Potential change with phase change**
- 5 Reference States
 - Reference State for Helmholtz free energy
 - Reference State for Chemical Potential
- 6 Fundamental Equations with Multiphysics Effects
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Chemical Potential change with phase change

- At $T = T_e$, α and β phases are under equilibrium, for α phase,

$$h^\alpha(T, p_0) = h^{\circ, \alpha}(T_0, p_0) + \int_{T_0}^T c_p^\alpha dT$$

$$s^\alpha(T, p_0) = s^{\circ, \alpha}(T_0, p_0) + \int_{T_0}^T \frac{c_p^\alpha}{T} dT$$

$$\mu^\alpha(T, p_0) = \mu^{\circ, \alpha}(T_0, p_0) + \int_{T_0}^T c_p^\alpha dT - T \int_{T_0}^T \frac{c_p^\alpha}{T} dT$$

- For β phase,

$$h^\beta(T, p_0) = h^\alpha(T_e, p_0) + \Delta h_e^\circ(T_e, p_0) + \int_{T_e}^T c_p^\beta dT$$

$$s^\beta(T, p_0) = s^\alpha(T_e, p_0) + \Delta s_e^\circ(T_e, p_0) + \int_{T_e}^T \frac{c_p^\beta}{T} dT$$

Chemical Potential change with phase change

- The chemical potential of β phase,

$$\begin{aligned}\mu^\beta(T, p_0) &= h^\beta(T, p_0) - T s^\beta(T, p_0) \\ &= h^\alpha(T_e, p_0) - T s^\alpha(T_e, p_0) + \Delta h_e^\circ(T_e, p_0) - T \Delta s_e^\circ(T_e, p_0) \\ &\quad + \int_{T_e}^T c_p^\beta dT - T \int_{T_e}^T \frac{c_p^\beta}{T} dT\end{aligned}$$

- By rearrangement,

$$\mu^\beta(T, p_0) = \mu^\beta(T_e, p_0) + \int_{T_e}^T [c_p^\beta - s^\beta(T_e, p_0)] dT - T \int_{T_e}^T \frac{c_p^\beta}{T} dT$$

- At (T_e, p_0) ,

$$\begin{aligned}\mu^\beta(T_e, p_0) &= \mu^\alpha(T_e, p_0) \\ &= h^{\circ, \alpha}(T_0, p_0) - T_e s^{\circ, \alpha}(T_0, p_0) + \int_{T_0}^{T_e} \frac{(T - T_e) c_p^\alpha}{T} dT\end{aligned}$$

Table of Contents

- 1 Choice of Fundamental Equation of Thermodynamics
- 2 Molar Helmholtz Free energy change
- 3 Chemical Potential change
- 4 Chemical Potential change with phase change
- 5 Reference States**
 - Reference State for Helmholtz free energy
 - Reference State for Chemical Potential
- 6 Fundamental Equations with Multiphysics Effects
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Reference State for Helmholtz free energy

- The reference state for molar Helmholtz free energy has almost always been associated with the state at 0 K,

$$f(T = 0 \text{ K}, v) = u(T = 0 \text{ K}, v)$$

- By third law of thermodynamics,

$$\lim_{T \rightarrow 0 \text{ K}} s(T, v_0) \rightarrow 0 \quad \lim_{T \rightarrow 0 \text{ K}} c_v(T) \rightarrow 0$$

$$\lim_{T \rightarrow 0 \text{ K}} c_p(T) \rightarrow 0 \quad \lim_{T \rightarrow 0 \text{ K}} \alpha(T) \rightarrow 0$$

- The Helmholtz free energy is

$$\begin{aligned} f(T, v) = & u(v) + \int_{0 \text{ K}}^T c_v dT - T \int_{0 \text{ K}}^T c_v d(\ln T) - p_0(v - v_0) \\ & - \int_{v_0}^v B_T dv + v \int_{v_0}^v B_T d(\ln v) \end{aligned}$$



Reference State for Chemical Potential

- In general, the enthalpy of the pure element at 1 bar and $T = 298 \text{ K}$,

$$h^\circ(298 \text{ K}, 1 \text{ bar}) = 0$$



Table of Contents

- 1 Choice of Fundamental Equation of Thermodynamics
- 2 Molar Helmholtz Free energy change
- 3 Chemical Potential change
- 4 Chemical Potential change with phase change
- 5 Reference States
 - Reference State for Helmholtz free energy
 - Reference State for Chemical Potential
- 6 Fundamental Equations with Multiphysics Effects**
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- The differential internal energy of the system

$$dU = TdS - pdV + V_0(\sigma_{ij}^d d\varepsilon_{ij}^d + E_i dD_i) + \phi dq + \gamma dA + \mu dN_i$$

where σ_{ij}^d and ε_{ij}^d are deviatoric components of stress and strain and γ is the interface energy.

- The differential Helmholtz and Gibbs free energies are

$$dF = -SdT - pdV + V_0(\sigma_{ij}^d d\varepsilon_{ij}^d + E_i dD_i) + \phi dq + \gamma dA + \mu dN_i$$

$$dG = -SdT + Vdp + V_0(\varepsilon_{ij}^d d\sigma_{ij}^d + E_i dD_i) + \phi dq + \gamma dA + \mu dN_i$$

- The volume densities at constant T and p are

$$du_v = Tds_v - pd\varepsilon + \sigma_{ij}^d d\varepsilon_{ij}^d + E_i dD_i + \phi d\rho_q + \gamma da_v + \mu_i dc_i$$

$$df_v = -pd\varepsilon + \sigma_{ij}^d d\varepsilon_{ij}^d + E_i dD_i + \phi d\rho_q + \gamma da_v + \mu_i dc_i$$

$$dg_v = -\varepsilon_{ij}^d d\sigma_{ij}^d + E_i dD_i + \phi d\rho_q + \gamma da_v + \mu_i dc_i$$

Electrochemical Potential

- At constant T and V and constant interface area and displacement,

$$\begin{aligned}dF &= \phi dq + \mu_i dN_i = \left(\phi \frac{dq}{dN_i} + \mu_i \right) dN_i \\ &= (\phi z_i e N_A + \mu_i) dN_i = (\phi z_i \mathcal{F} + \mu_i) dN_i\end{aligned}$$

where z_i is the valence of the species i , e is the elementary charge, N_A is Avogadro's number.

- Let the electrochemical potential of species i is

$$\tilde{\mu}_i = \mu_i + z_i \mathcal{F} \phi$$



Chemical Potential Including Surface

- At constant T and p ,

$$dG = \gamma dA + \mu dN$$

Proceed to

$$dG = \gamma \left(\frac{\partial A}{\partial V} \right) \left(\frac{\partial V}{\partial N} \right) dN + \mu dN = \gamma v \left(\frac{\partial A}{\partial V} \right) dN + \mu dN$$

it becomes

$$dG = \left[\gamma v \left(\frac{\partial A}{\partial V} \right) + \mu \right] dN$$

Let

$$\mu_r = \gamma v \left(\frac{\partial A}{\partial V} \right) + \mu$$

For solid particle r ,

$$\mu_r = \gamma v \left(\frac{\partial A / \partial r}{\partial V / \partial r} \right)_{T,p} + \mu$$



Chemical Potential Including Surface

- Assume the particle is spherical,

$$A = 4\pi r^2 \quad V = \frac{4\pi}{3} r^3$$

therefore,

$$\mu_r = \frac{2\gamma v}{r} + \mu$$

- Since $r \rightarrow \infty$,

$$\mu = \mu_\infty$$

